

Dedicated to Professor Ioan Gottlieb's  
80<sup>th</sup> Anniversary

## THE ISSUE OF PRIMITIVE INDEFINITE-INTEGRAL IN THE THEORY OF FRACTAL SPACE-TIME

I. GOTTLIEB<sup>1,4</sup>, C. MOCIUTCHI<sup>1</sup>, G. CIOBANU<sup>1</sup>, M. GARTU<sup>2</sup>, V. GARTU<sup>2</sup>,  
M. COLOTIN<sup>1</sup>, M. AGOP<sup>3</sup>

<sup>1</sup> Faculty of Physics, "A.I. Cuza" University, Blvd. Carol I, No. 11, Iasi 700506 Romania,  
e-mail: gottlieb1929@yahoo.com

<sup>2</sup> Department of Mathematics, University of Bacau, Calea Marasesti No. 157, Bacau, Romania

<sup>3</sup> Department of Physics, Technical "Gh. Asachi" University, Blvd. Mangeron, Iasi, 700029, Romania

<sup>4</sup> Corresponding Author: Faculty of Physics, "A.I. Cuza" University, Blvd. Carol I, No. 11, Iasi  
700506 Romania

(Received June 4, 2009)

*Abstract.* Using an interpolation method, the issue of primitive - indefinite integral in the theory of fractal space-time is established. In this context, the von Koch curve case is analyzed. Moreover, a correspondence with Cantorian space-time is determined.

*Key words:* fractal space-time, Cantor set, indefinite-integral.

### 1. INTRODUCTION

Let us consider a fractal curve  $F$  which is self affine [1, 2]. It has no derivative in none of its points. The necessity of connecting the fractal theory with the results [3-6], where the derivatives play an essential role, leads us to defining of the left and the right derivatives of a certain point of the curve. Their average values are connected with the usual derivative, while their semi-difference characterizes the fractality (for other details see [2–6]).

However, there are a number of problems in the field of Physics, where the notion of integral is also required [6-12]. This aspect will be analyzed in the followings. First, some specifications must be considered:

a) A curve  $I$ , which is the primitive of  $F$ , if there is at least in an approximation, must be continuous and derivable, hence of class  $C_1$ , since its derivative in any point should render a point of  $F$ , which is well-determined;

b) The analytical expression of  $I$  in the form  $y = y(x)$  is difficult to be established. Instead of this, we can state a parameterization

$$x = x(t); y = y(t); \quad (1)$$

with  $t \in (t_A, t_B)$  to describe a curve  $F$  between the points

$$A(x(t_A), y(t_A)) \in F \text{ and } B(x(t_B), y(t_B)) \in F. \quad (2)$$

As a rule, (1) could be achieved by choosing a value  $t' \in (t_A, t_B)$  for an intermediary point on  $F$  between  $A$  and  $B$  and by continuing this course endlessly;

c) In the next step we shall choose an approximation of (1). For  $n + 1$  values of  $t$ , denoted by  $t_0, t_1, \dots, t_i, \dots, t_n$ , that satisfy

$$t_A = t_0 < t_1 < \dots < t_n = t_B, \quad (3)$$

let us assume that the values of  $x(t_i) = x_i$  and  $y(t_i) = y_i$  are known. Then, we can build up the functions

$$x = X(t) \text{ and } y = Y(t) \quad (4)$$

so that

$$X(t_i) = x(t_i) \text{ and } Y(t_i) = y(t_i). \quad (5)$$

The curves (4) are interpolation functions on the known values of (1) in the points (3).

## 2. EXAMPLES OF INTERPOLATION

There are many possibilities to obtain the interpolation functions (5) – for details see [7], - which are continuous, hence of class  $C_0$ , as well as  $F$ . Three of these possibilities are presented as follows:

a) One of them is to connect the neighboring points by a segment. Between the points with the coordinates  $x(t_{i-1})$ ,  $x(t_i)$  and  $x(t_{i+1})$ ,  $0 < i < n$  we shall have two segments whose equations are

$$\left\{ \begin{array}{l} x(t) = x(t_{i-1}) + \frac{x(t_i) - x(t_{i-1})}{t_i - t_{i-1}} t, \quad \text{with} \quad t_{i-1} < t < t_i \\ x(t) = x(t_i) + \frac{x(t_{i+1}) - x(t_i)}{t_{i+1} - t_i} t, \quad \text{with} \quad t_i < t < t_{i+1} \end{array} \right.$$

with the primitives

$$\left\{ \begin{array}{l} \int x(t) dt = x(t_{i-1})t + \frac{x(t_i) - x(t_{i-1})}{t_i - t_{i-1}} \frac{t^2}{2} + C_{i-1}, \quad \text{with } t_{i-1} < t < t_i \\ \int x(t) dt = x(t_i)t + \frac{x(t_{i+1}) - x(t_i)}{t_{i+1} - t_i} \frac{t^2}{2} + C_i, \quad \text{with } t_i < t < t_{i+1}. \end{array} \right. \quad (6)$$

By imposing the continuity of these primitives, one can determine the connection between the integration constants  $C_{i-1}$  and  $C_i$ . This fact provides the existence of the derivative for  $t_i$  which coincides with the value  $x(t_i)$  of the interpolation function.

Perhaps this is the closest way to the fractal representation.

b) Another possibility is to choose the Lagrange interpolation polynomials of the order  $n$ . To obtain such an interpolation polynomial, let us denote

$$\left\{ \begin{array}{l} P_n(t) = a_0 + a_1t + \dots + a_nt^n \\ Q_n(t) = b_0 + b_1t + \dots + b_nt^n. \end{array} \right. \quad (7a,b)$$

The conditions (5) require

$$P_n(t_i) = x(t_i) \text{ and } Q_n(t_i) = y(t_i), \text{ with } i = 0, \dots, n. \quad (8)$$

Under such circumstances, let us expand the  $P_n(t)$  polynomial. We have

$$\left\{ \begin{array}{l} a_0 + a_1t_0 + \dots + a_nt_0^n = x(t_0) \\ \dots \quad \dots \quad \dots \\ b_0 + b_1t_n + \dots + b_nt_n^n = x(t_n), \end{array} \right. \quad (9)$$

which forms a system of  $(n+1)$  equations for the  $(n+1)$  coefficients  $a_0, \dots, a_n$  of the required polynomial. By adding (7a) to the system (9), we obtain  $(n+2)$  equations for these coefficients. The compatibility condition gives

$$\begin{vmatrix} 1 & t_0 & \dots & t_0^n & x(t_0) \\ \dots & \dots & \dots & \dots & \dots \\ 1 & t_n & \dots & t_n^n & x(t_n) \\ 1 & t & \dots & t^n & P_n(t) \end{vmatrix} = 0,$$

*i.e.*

$$\begin{vmatrix} 1 & t_0 & \dots & t_0^n & x(t_0) \\ \dots & \dots & \dots & \dots & \dots \\ 1 & t_n & \dots & t_n^n & x(t_n) \\ 0 & 0 & \dots & 0 & P_n(t) \end{vmatrix} + \begin{vmatrix} 1 & t_0 & \dots & t_0^n & x(t_0) \\ \dots & \dots & \dots & \dots & \dots \\ 1 & t_n & \dots & t_n^n & x(t_n) \\ 1 & t & \dots & t^n & 0 \end{vmatrix} = 0,$$

and from here

$$P_n(t) = - \frac{\begin{vmatrix} 1 & t_0 & \dots & t_0^n & x(t_0) \\ \dots & \dots & \dots & \dots & \dots \\ 1 & t_n & \dots & t_n^n & x(t_n) \\ 1 & t & \dots & t^n & 0 \end{vmatrix}}{\begin{vmatrix} 1 & t_0 & \dots & t_0^n \\ \dots & \dots & \dots & \dots \\ 1 & t_n & \dots & t_n^n \end{vmatrix}}. \quad (10)$$

The  $Q_n(t)$  polynomial is also obtained in a similar manner.

c) We can also accomplish the interpolation by the “Fourier sums”, *i.e.* by a part of a Fourier’s series, named trigonometrical polynomials, such as

$$T(t) = a_0 + \sum_{k=1}^m (a_k \cos kt + b_k \sin kt). \quad (11)$$

By taking into account that the number of the coefficients to be determined is  $2m+1$ , we shall require the same number of points in which the function is known. The determination methods for  $T(t)$  are similar to those used in b). Thus, by imposing (11), we obtain

$$a_0 + \sum_{k=1}^m (a_k \cos kt + b_k \sin kt) = x(t_i), \quad i = \overline{0, 2m}. \quad (12)$$

The compatibility of the equations (12) and (11) for the unknown values  $a_k$ ,  $b_k$  ( $k = \overline{0, 2m}$ ) requires

$$\begin{vmatrix} 1 & \cos t_0 & \sin t_0 & \dots & \cos mt_0 & \sin mt_0 & x(t_0) \\ \dots & \dots & \dots & \dots & \dots & \dots & \dots \\ 1 & \cos t_{2m} & \sin t_{2m} & \dots & \cos mt_{2m} & \sin mt_{2m} & x(t_{2m}) \\ 1 & \cos t & \sin t & \dots & \cos mt & \sin mt & T(t) \end{vmatrix} = 0. \quad (13)$$

Similarly, as in the case b),  $T(t)$  is the ratio of two determinants, and represents the interpolation function of  $x(t)$  by trigonometrical polynomials.

d) Instead of the trigonometrical functions, one can have orthogonal functions such as the eigenfunctions from the quantum mechanics in the case of a discrete spectrum.

### 2.1. The Case of the Von Koch Curve

As an example, let us consider the von Koch curve for the second iteration (Fig. 1). Since the intervals  $p_i - p_{i-1}$  ( $i = \overline{1,16}$ ) are equal, we shall consider the parameter  $t$  divided in a similar number of coordinates of equal parts. To each  $t_i$  we have an  $x_i$  and an  $y_i$ . These are the coordinates of the  $p_i$  point. The corresponding values are systematized in Table 1. These values of  $t$  are given exactly, the ones of  $x_i$  similarly, with the condition that the underlined decimal is infinitely repeated, while the ones of  $y_i$  are given with three decimal precision.

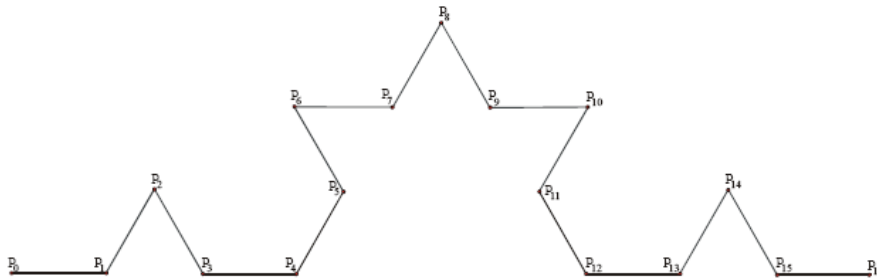


Fig. 1 – The von Koch curve for the second iteration.

Table 1

The values of coordinates  $x_i(t)$  and  $y_i(t)$  of the points  $p_i$

	$P_0$	$P_1$	$P_2$	$P_3$	$P_4$	$P_5$	$P_6$	$P_7$	$P_8$	$P_9$	$P_{10}$	$P_{11}$	$P_{12}$	$P_{13}$	$P_{14}$	$P_{15}$	$P_{16}$
$t$	0	0,0625	0,125	0,1875	0,25	0,3125	0,375	0,4375	0,5	0,5625	0,625	0,6875	0,75	0,8125	0,875	0,9375	1,0
$x$	0	0,111	0,166	0,222	0,333	0,388	0,333	0,444	0,5	0,555	0,666	0,611	0,666	0,777	0,833	0,888	1,0
$y$	0	0	0,096	0	0	0,086	0,192	0,192	0,289	0,192	0,193	0,096	0	0	0,096	0	0

We shall exemplify the above mentioned facts for the case of the von Koch curve, approximated by segments (Fig. 2) and we shall refer to the first 5 segments defined by the first 6 values of  $t$ , i.e. for

$$t_0 = 0; \quad t_1 = \alpha; \quad t_2 = 2\alpha; \quad t_3 = 3\alpha; \quad t_4 = 4\alpha, \quad \text{with } \alpha = (0.25)^2.$$

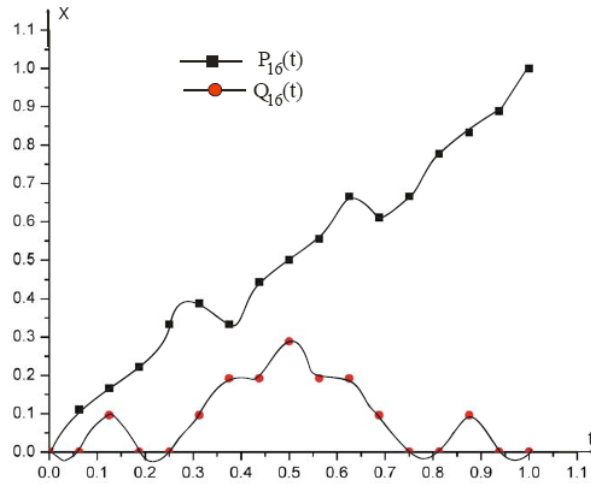


Fig. 2 – The segmented curves for the interpolation of  $x(t)$  and  $y(t)$ .

Thus, we obtain (Fig. 3):

$$x(t) = \begin{cases} \frac{0.111}{\alpha} t & \dots & \dots & 0 < t < \alpha \\ 0.111 + \frac{0.055}{\alpha} (t - \alpha) & \dots & \dots & \alpha < t < 2\alpha \\ 0.166 + \frac{0.056}{\alpha} (t - 2\alpha) & \dots & \dots & 2\alpha < t < 3\alpha \\ 0.222 + \frac{0.111}{\alpha} (t - 3\alpha) & \dots & \dots & 3\alpha < t < 4\alpha \\ 0.333 + \frac{0.055}{\alpha} (t - 4\alpha) & \dots & \dots & 4\alpha < t < 5\alpha, \end{cases}$$

and

$$y(t) = \begin{cases} 0 & \dots & \dots & 0 < t < \alpha \\ \frac{0.096}{\alpha} (t - \alpha) & \dots & \dots & \alpha < t < 2\alpha \\ 0.096 - \frac{0.096}{\alpha} (t - 2\alpha) & \dots & \dots & 2\alpha < t < 3\alpha \\ 0 & \dots & \dots & 3\alpha < t < 4\alpha \\ \frac{0.096}{\alpha} (t - 4\alpha) & \dots & \dots & 4\alpha < t < 5\alpha. \end{cases}$$

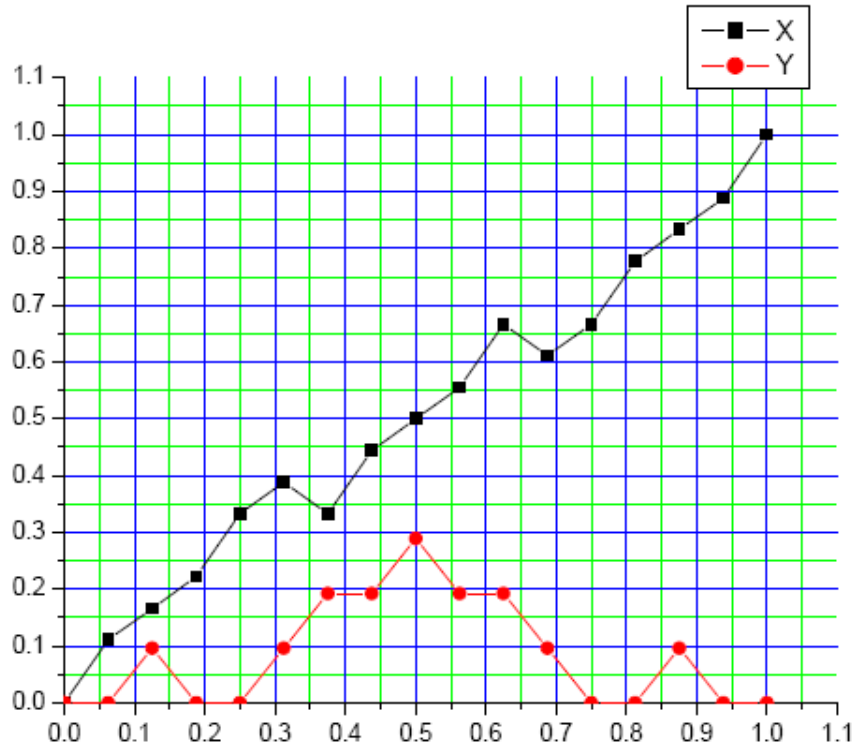


Fig. 3 – The approximation interpolation of  $x(t)$  and  $y(t)$  by the  $C^\infty$  class functions (polynomials, trigonometrical polynomials, etc.).

From here, by integration, we obtain (Fig. 4):

$$\int_0^t x(t) dt = \begin{cases} \frac{0.111}{2\alpha} t^2 & 0 < t < \alpha \\ \frac{0.111}{2} \alpha + 0.111(t - \alpha) + \frac{0.055}{2\alpha} (t - \alpha)^2 & \alpha < t < 2\alpha \\ \frac{0.388}{2} \alpha + 0.166(t - 2\alpha) + \frac{0.056}{2\alpha} (t - 2\alpha)^2 & 2\alpha < t < 3\alpha \\ \frac{0.776}{2} \alpha + 0.222(t - 3\alpha) + \frac{0.111}{2\alpha} (t - 3\alpha)^2 & 3\alpha < t < 4\alpha \\ \frac{0.776}{2} \alpha + 0.333(t - 4\alpha) + \frac{0.055}{2\alpha} (t - 4\alpha)^2 & 4\alpha < t < 5\alpha, \end{cases}$$

and

$$\int_0^t y(t) dt = \begin{cases} 0 & 0 < t < \alpha \\ \frac{0.096}{2\alpha}(t - \alpha)^2 & \alpha < t < 2\alpha \\ 0.096(t - 2\alpha) - \frac{0.096}{2\alpha}(t - 2\alpha)^2 & 2\alpha < t < 3\alpha \\ \frac{0.096}{2}\alpha & 3\alpha < t < 4\alpha \\ \frac{0.096}{2}\alpha + \frac{0.096}{2\alpha}(t - 4\alpha)^2 & 4\alpha < t < 5\alpha. \end{cases}$$

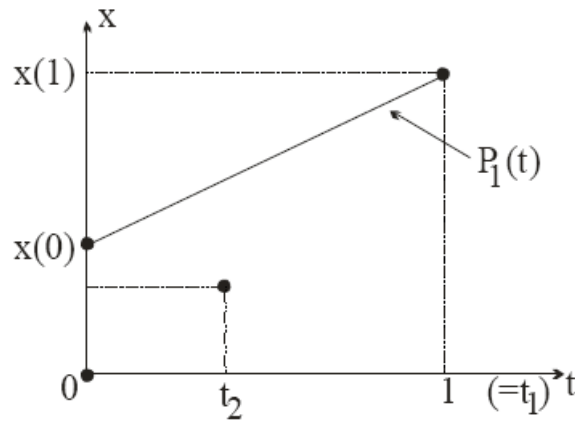


Fig. 4 – Interpolating polynomial of the first order.

The other mentioned approximations lead to similar results; they are more laborious, but simpler to be programmed on computer.

## 2.2. The correspondence with Cantorian space-time

As it was also shown in paragraph 1.2, we summarize at  $x(t)$ . This function can be approximated by polynomials. We start by the interpolating polynomial of the first order. For  $t=0$  and  $t=1$  we have

$$P_1(0) = x(0) \quad \text{and} \quad P_1(1) = x(1) \quad (14)$$

$P_1(t)$  being a line (a line segment) – Fig. 4.



Now, we add an intermediary point (it must not be in the middle!)  $t_2$ , in which the function takes the value  $x(t_2)$ . The polynomial of the second order (a parabola) will pass through the three points,

$$P_2(0) = x(0); \quad P_2(t_2) = x(t_2); \quad P_2(1) = x(t_1). \quad (15)$$

Then, we add another point, keeping the previous ones, thus obtaining polynomials with an extra degree. Thus we find  $P_n(t)$  for which

$$P_n(0) = x(0); \quad P_n(t_1) = x(t_1); \quad P_n(t_2) = x(t_2); \quad \dots; \quad P_n(t_{n-1}) = x(t_{n-1}) \quad (16)$$

having the form

$$P_n = a_0 t^n + a_1 t^{n-1} + \dots + a_n. \quad (17)$$

**Theorem.** *We assert that the sequence  $\{P_n\}$  is convergent and*

$$\lim_{n \rightarrow \infty} P_n(t) = x(t) \quad (18)$$

*if we ascertain the fact that the values sequence, all different, for  $t$  should be compact on the entire interval  $(0, 1)$ . (The set  $B$  is compact in the set  $A$  if the closure of  $B$  is equal with  $A$ ).*

Actually, let us consider  $t = t_k$  and the sequential values of these interpolating polynomials in this point,

$$P_1(t_k), \dots, P_k(t_k), P_{k+1}(t_k, \dots). \quad (19)$$

From the way these polynomials are built, it results

$$P_k(t_k) = P_{k+1}(t_k) = \dots = P_n(t_k) \quad \text{with} \quad n \geq k. \quad (20)$$

Therefore, this will also be the limit of the sequence in  $t = t_k$ .

If we ascertained the fact that the sequence  $\{t_k\}$  is compact in the interval  $[0, 1]$ , then for the values  $t \notin \{t_k\}$  we find an approximation as good as possible for  $t$ .

One of the Lebesgue theorems [2-5] asserts that a continuous curve, which is not differentiable in any point, has an infinite length. Since our curve is almost all over fractal self-similar, it satisfies the condition of the Lebesgue theorems. Between the  $t'$  and  $t''$  values, we can approximate the length of the fractal curve with

$$\int_{t'}^{t''} \sqrt{(dP_n)^2 + (dQ_n)^2} = \int_{t'}^{t''} \sqrt{\left(\frac{dP_n}{dt}\right)^2 + \left(\frac{dQ_n}{dt}\right)^2} dt. \quad (21)$$

When  $n \rightarrow \infty$ , the limit of these integrals has to be infinite. Thus, either  $dP_n/dt$  or  $dQ_n/dt$  (both, according to our opinion) should diverge. Consequently, either the curve  $x(t)$  or  $y(t)$  are fractal (or both!).

However, let us consider for the moment that  $n$  is no matter how great but finite and, by an adequate normalization, let us suppose

$$\left(\frac{dP_n}{dt}\right)^2 + \left(\frac{dQ_n}{dt}\right)^2 = X^2 + Y^2 = 1. \quad (22)$$

Then, the equation (22) may be interpreted as absolute in the  $X$  and  $Y$  coordinates and we obtain the metric [13]

$$ds^2 = -\frac{dh d\bar{h}}{(h - \bar{h})^2}, \quad (23)$$

for

$$h = \frac{Y \pm i\sqrt{1 - X^2 - Y^2}}{1 - X}. \quad (24)$$

Since the metric (23) is invariant to the  $SL(2R)$  group, the restriction to the quantum space-time implies the subgroup  $SL(2Z)$  that is the transformations [13]:

$$h' = \frac{ah + b}{ch + d}, \quad \bar{h}' = \frac{a\bar{h} + b}{c\bar{h} + d}, \quad a, b, c, d \in Z \quad ad - bc = \pm 1 \quad (25a-d)$$

Now, let us consider a transformation element of the matrix form

$$M = \begin{pmatrix} a & b \\ c & d \end{pmatrix}.$$

Under special conditions,  $c = 1$ ,  $d = 0$ , we can calculate the eigenvalues by means of the irrational roots [13]

$$x_1^{(n)} = (-\phi)^n, \quad x_2^{(n)} = (1/\phi)^n,$$

where  $\phi = (\sqrt{5} - 1)/2$  is the golden mean of Cantorian (El Naschie's)  $\varepsilon^{(\infty)}$  space-time theory [14-21]. This clearly shows the connection between the geometry induced by the previous transformation, KAM theorem, the VAK and El Naschie's  $\varepsilon^{(\infty)}$  theory [15-21].

### 3. CONCLUSIONS

The aim of the present paper was to establish a framework of using the integro-differential equations in the fractal theory. The main conclusions are as follows: i) using an interpolation method, the issue of primitive - indefinite integral in the theory of fractal space-time is established; ii) some examples of interpolation are given (interpolation by segments, Lagrange interpolation polynomials, interpolation by the Fourier's series); iii) the von Koch curve case is analyzed; iv) a correspondence with El Naschie's  $\varepsilon^{(\infty)}$  space-time is determined.

*Acknowledgement.* We wish to thank to Professor M.S. El Naschie for helpful discussion on the manuscript of this paper.

### REFERENCES

1. B. Mandelbrot, *The fractal geometry of nature*, Freeman, San Francisco, 1982.
2. L. Nottale, *Fractal Space-Time and Microphysics: Towards a Theory of Scale Relativity*, World Scientific, Singapore, 1993.
3. J. Feder and A. Aharony, *Fractals in Physics*, North Holland, Amsterdam, 1990.
4. J.F. Gouyet, *Physique et Structures Fractals*, Masson, Paris, 1992.
5. M. Barnsley, *Fractals Everywhere. Deterministic Fractal Geometry* Academic Press Boston, 1988.
6. M.S. El Naschie, O.E. Röslér, I. Prigogine, *Quantum Mechanics, Diffusion and Chaotic Fractals*, Elsevier, Oxford, 1995.
7. Elena Popovici, *The average theorems from mathematical analyses and there connection with the interpolation theory*, Edit. Dacia, Cluj, Romania, 1972.
8. J. Argyris, C. Ciobotariu, and G. Mattutis, *Chaos, Solitons and Fractals*, **12**, 1 (2001).
9. J. Argyris, C. Marin, and C. Ciobotariu, *Physics of Gravitation and the Universe*, Spiru Haret Publishing House, Iasi, 2003.
10. R.P.Feynman and A.R. Hibbs, *Quantum mechanics and Path Integrals*, McGraw-Hill, New-York, 1965.
11. L.F. Abbott and M.B. Wise, *Am. J.Phys.*, **49**, 37 (1981).
12. E. Campesina-Romeo, J.C. D'Oliovo and M. Socolovsky, *Phys.Lett.*, **89A(7)**, 321 (1982).
13. Gottlieb I., Agop M., Jarcau M., *Chaos, Solitons and Fractals*, **19**, 705 (2004).
14. Weibel P., Ord. G., Rössler O. (Editors), *Space-time physics and Fractality*, Festschrift in honor of Mohamed El Naschie, Vienna, New York, Springer, 2005.
15. El Naschie MS, *Chaos, Solitons and Fractals*, **22**, 495 (2004).
16. El Naschie MS, *Chaos, Solitons and Fractals*, **22**, 1 (2004).
17. El Naschie MS, *Chaos, Solitons and Fractals*, **8**, 131 (1997).
18. El Naschie MS, *Chaos, Solitons and Fractals*, **8**, 1865 (1997).
19. El Naschie MS, *Int. J. of Nonlinear Sciences and Numerical Simulation*, **8(3)**, 445 (2007).
20. Mukhamedov A, *Chaos, Solitons and Fractals* **33** 1 (2007).
21. Mukhamedov A, *Chaos, Solitons and Fractals* **33**, 717 (2007).